

10. Electricity and Magnetism (Spring 2006)

An insulated, spherical, conducting shell of radius  $a$  is in a uniform electric field  $E_0$ . If the sphere is cut into two hemispheres by a plane perpendicular to the field, find the force required to prevent the hemispheres from separating

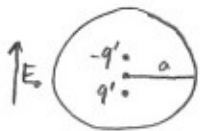
(a) if the shell is uncharged;

This is Jackson 2.9

(b) if the total charge on the shell is  $Q$ .

$-q_0 (z=R)$

a. The strategy is:



$q_0 (z=-R)$

Method of Images:

$z' = \frac{a^2}{R} \quad q' = -\frac{a}{R} q$

1) Imagine the  $E_0$  field is created by two point charges a large distance  $R$  away and use the method of images to find the potential outside the sphere

2) Find  $\sigma$  using  $\sigma = -\epsilon_0 \frac{\partial V}{\partial n} |_{r=a}$

3) Find the force using  $F_z = 2 \int_{\text{hemi}} E_z dq$

$= 2 \int_{\text{hemi}} \left( \frac{E}{2\epsilon_0} \cos(\theta) \right) \sigma dA$

$= 2 \int_0^{2\pi} \int_0^{\pi/2} \frac{\sigma^2}{2\epsilon_0} \cos(\theta) a^2 \sin(\theta) d\theta d\phi$

where the factor of 2 comes from having to push both hemispheres in order to hold them together.

1)  $E_0 = 2 \frac{1}{4\pi\epsilon_0} \frac{q}{R^2} \Rightarrow q = 2\pi\epsilon_0 R^2 E_0$

Recall  $\vec{p} \equiv \int \vec{r}' \rho(\vec{r}') d^3r' \Rightarrow \vec{p}$  points toward positive charge, unlike fields

$\vec{p} = |2z'| \cdot |q'| \hat{z} = 2 \frac{a^3}{R^2} q \hat{z} = 4\pi\epsilon_0 a^3 E_0 \hat{z}$

Remember that  $V_{\text{dip}} = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2}$  so

$V_{\text{dip}} = \frac{1}{4\pi\epsilon_0} \frac{4\pi\epsilon_0 a^3 E_0}{r^2} \cos(\theta) = \frac{a^3}{r^2} E_0 \cos(\theta)$

$V_{\text{field}} = -\int \vec{E} \cdot d\vec{r} = -E_0 r \cos(\theta)$

$V = V_{\text{dip}} + V_{\text{field}} = -\left(r - \frac{a^3}{r^2}\right) E_0 \cos(\theta)$

2)  $\sigma = -\epsilon_0 \frac{\partial V}{\partial n} |_{r=a} = -\epsilon_0 \frac{\partial V}{\partial r} |_{r=a} = \epsilon_0 \left(1 + 2\frac{a^3}{r^3}\right) E_0 \cos(\theta) |_{r=a}$

$= 3\epsilon_0 E_0 \cos(\theta)$  which is Jackson (2.15)

3)  $F_z = q \epsilon_0 E_0^2 a^2 (2\pi) \int_0^{\pi/2} \cos^3(\theta) \sin(\theta) d\theta$  Let  $x = \cos(\theta)$

$= 9\epsilon_0 E_0^2 a^2 (2\pi) \int_0^1 x^3 dx = \frac{9}{2} \pi \epsilon_0 E_0^2 a^2$

b. In part a, the surface was an equipotential, so an additional charge  $Q$  will spread out uniformly:  $\sigma \rightarrow 3\epsilon_0 E_0 \cos(\theta) + Q/4\pi a^2$

$F_z = \frac{2\pi a^2}{\epsilon_0} \int_0^{\pi/2} \left[ 9\epsilon_0 E_0 \cos^2(\theta) + 2 \frac{3\epsilon_0 E_0 Q}{4\pi a^2} \cos(\theta) + \frac{Q^2}{16\pi^2 a^4} \right] \cos(\theta) \sin(\theta) d\theta$

$= \frac{9}{2} \pi \epsilon_0 E_0^2 a^2 + 3E_0 Q \int_0^{\pi/2} \cos^2(\theta) \sin(\theta) d\theta + \frac{Q^2}{8\pi a^2 \epsilon_0} \int_0^{\pi/2} \cos(\theta) \sin(\theta) d\theta$

$= \frac{9}{2} \pi \epsilon_0 E_0^2 a^2 + E_0 Q + \frac{Q^2}{16\pi a^2 \epsilon_0}$

But the middle term is the  $Q$  charge interacting with the field, which is in the same direction for both hemispheres, so

$F = \frac{9}{2} \pi \epsilon_0 E_0^2 a^2 + \frac{Q^2}{16\pi a^2 \epsilon_0}$